Fast quantum integer multiplication with almost no ancillas

Gregory D. Kahanamoku-Meyer July 14, 2023

Motivation

Arithmetic on quantum computers: why do we care?

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OPEN

Classically verifiable quantum advantage from a computational Bell test

Gregory D. Kahanamoku-Meyer^{®1™}, Soonwon Choi¹, Umesh V. Vazirani^{2™} and Norman Y. Yao^{®1™}

Existing experimental demonstrations of quantum computational advantage have had the limitation that verifying the correctness of the quantum device requires exponentially costly classical computations. Here we propose and analyse an interactive protocol for demonstrating quantum computational advantage, which is efficiently classically verifiable. Our protocol relies on a class of cryptographic tools called trapdoor claw-free functions. Although this type of function has been applied to quantum advantage protocols before, our protocol employs a surprising connection to Bell's inequality to avoid the need for a demanding cryptographic property called the adaptive hardroot bit, while maintaining essentially no increase in the quantum circuit complexity and no extra assumptions. Leveraging the relaxed cryptographic requirements of the protocol, we present two trapdoor claw-free function constructions, based on Rabin's function and the Diffied-fellman problem, which have not been used in this context before. We also

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A Cryptographic Test of Quantumness and Certifiable Randomness from a Single Quantum Device

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URMILA MAHADEV, California Institute of Technology, USA

UMESH VAZIRANI, UC Berkeley, USA

THOMAS VIDICK, California Institute of Technology, USA

We consider a new model for the testing of untrusted quantum devices, consisting of a single polynomial time bounded quantum device interacting with a classical polynomial time verifier. In this model, we propose solutions to two tasks—a protocol for efficient classical verification that the untrusted device is "truly quantum" and a protocol for producing certifiable randomness from a single untrusted quantum device. Our solution relies on the existence of a new cryotographic primitive for constraining the power.

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POLYNOMIAL-TIME ALGORITHMS FOR PRIME FACTORIZATION AND DISCRETE LOGARITHMS ON A QUANTUM COMPUTER*

PETER W. SHOR

Abstract. A digital computer is generally believed to be an efficient universal computing device; that is, it is believed able to simulate any physical computing device with an increase in computation time by at most a polynomial factor. This may not be true when quantum mechanics is taken into consideration. This paper considers factoring integers and finding discrete logarithms, two problems which are generally thought to be hard on a classical computer and which have been used as the basis

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... with as few gates and qubits as possible.

Today's goal: implement the following unitaries

$$\mathcal{U}_{q \times q} \ket{x} \ket{y} \ket{w} = \ket{x} \ket{y} \ket{w + xy}$$

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Overview

1. Fast multiplication (few gates)

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- 2. Quantum multiplication (few qubits)

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- 1. Fast multiplication (few gates)
- 2. Quantum multiplication (few qubits)
- 3. Fast quantum multiplication (few gates + qubits)

The "schoolbook" method: $xy = \sum_{ij} (2^i x_i)(2^j y_j) = \sum_{ij} 2^{i+j} x_i y_j$

				1	1	0	1
			×	1	0	1	0
				1	0	1	0
		1	0	1	0		
	1	0	1	0			
1	0	0	0	0	0	1	0

5

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Running time: $\mathcal{O}(n^2)$ operations

Given two *n*-bit numbers *x* and *y*, what if we use base $b = 2^{n/2}$?

6

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$$\begin{array}{cccc}
 & x_1 & x_0 \\
 & \times & y_1 & y_0 \\
\hline
 & x_0 y_0 & \\
 & x_1 y_0 & \\
 & x_0 y_1 & \\
 & + & x_1 y_1 & \\
\hline
 & xy = x_1 y_1 b^2 + x_0 y_1 b + x_1 y_0 b + x_0 y_0
\end{array}$$

6

Given two *n*-bit numbers x and y, what if we use base $b = 2^{n/2}$?

$$\begin{array}{c|ccccc} & x_1 & x_0 \\ \times & y_1 & y_0 \\ \hline & & x_0 y_0 \\ \hline & & x_1 y_0 \\ & & x_0 y_1 \\ + & x_1 y_1 & & & \\ \end{array}$$

$$xy = x_1y_1b^2 + x_0y_1b + x_1y_0b + x_0y_0$$

Time remains $\mathcal{O}(n^2)$, because $4(n/2)^2 = n^2$

$$xy = x_1y_1b^2 + (x_0y_1 + x_1y_0)b + x_0y_0$$

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Observation:
$$x_0y_1 + x_1y_0 = (x_1 + x_0)(y_1 + y_0) - x_1y_1 - x_0y_0$$

Can compute xy with only three multiplications of size n/2:

- 1. x_1y_1
- 2. x_0y_0
- 3. $(x_1 + x_0)(y_1 + y_0)$

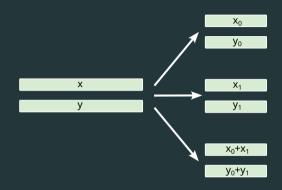
$$xy = x_1y_1b^2 + (x_0y_1 + x_1y_0)b + x_0y_0$$

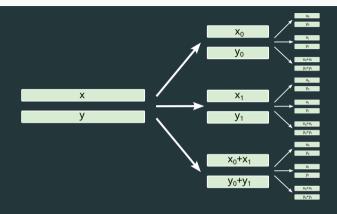
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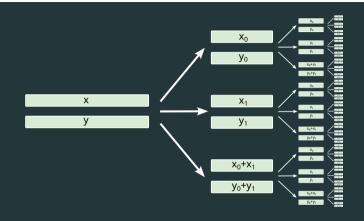
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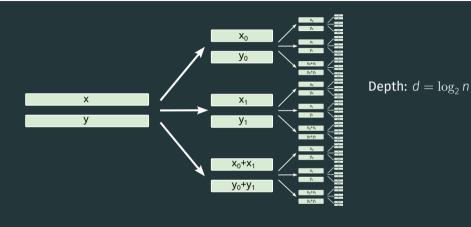
- 1. x_1y_1
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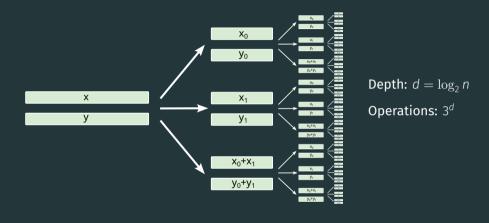
Computational cost:
$$3(n/2)^2 = \frac{3}{4}n^2 = \mathcal{O}(n^2)$$

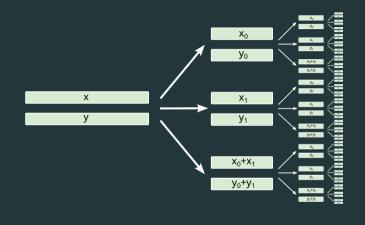








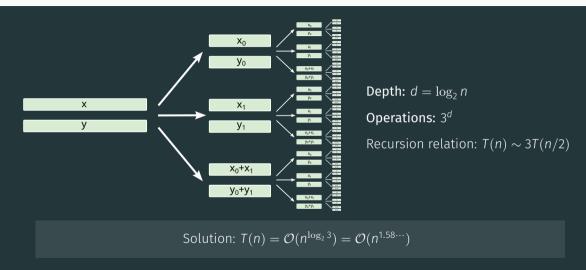




Depth: $d = \log_2 n$

Operations: 3^d

Recursion relation: $T(n) \sim 3T(n/2)$



Question: why don't we always do this, classically?

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GNU multiple-precision arithmetic library cutoff: 2176 bit numbers

Background: Toom-Cook multiplication

Break the n bit numbers into k chunks of n/k bits.

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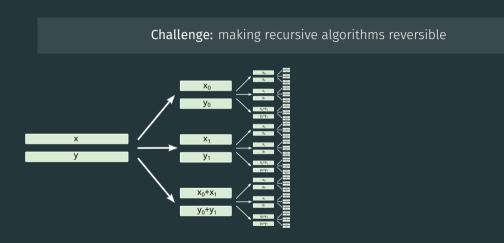
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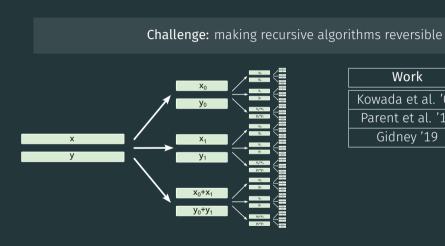
Algorithm	Gate count
Schoolbook	$\mathcal{O}(n^2)$
k = 2	$\mathcal{O}(n^{\log_2 3}) = \mathcal{O}(n^{1.58\cdots})$
k = 3	$\mathcal{O}(n^{\log_3 5}) = \mathcal{O}(n^{1.46\cdots})$
k = 4	$\mathcal{O}(n^{\log_4 7}) = \mathcal{O}(n^{1.40\cdots})$

Summary: fast multiplication

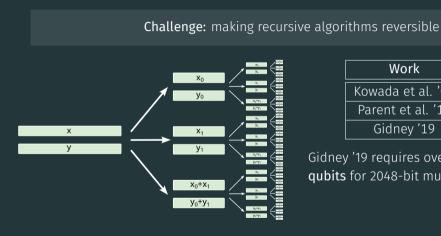
- "Standard" multiplication requires time $\mathcal{O}(n^2)$ operations
- $\boldsymbol{\cdot}$ Faster algorithms exist, but have large constant factors

Can these fast circuits be made quantum?



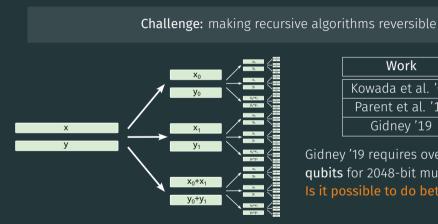


Work	Qubits
Kowada et al. '06	$\mathcal{O}(n^{1.58\cdots})$
Parent et al. '18	$O(n^{1.43})$
Gidney '19	$\mathcal{O}(n)$



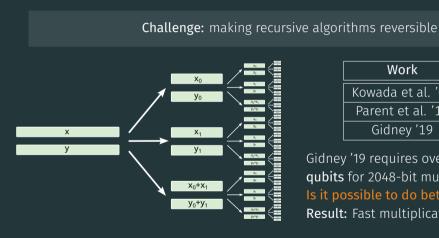
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Gidney '19 requires over 12,000 ancilla qubits for 2048-bit multiplication.



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Result: Fast multiplication using 1 ancilla

Quantum Fourier transform:

$$QFT |w\rangle = \sum_{z} \exp\left(\frac{2\pi i w z}{2^n}\right) |z\rangle$$

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$$|xy\rangle = QFT^{-1} \sum_{z} \exp\left(\frac{2\pi i x y z}{2^n}\right) |z\rangle$$

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How to implement
$$|x\rangle |y\rangle |0\rangle \rightarrow |x\rangle |y\rangle |xy\rangle$$
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1) Generate $|x\rangle|y\rangle\sum_{z}|z\rangle$, 2) apply a phase rotation of $\exp\left(\frac{2\pi i xyz}{2^n}\right)$, 3) apply QFT⁻¹

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$$xyz = \sum_{i,j,k} 2^i 2^j 2^k x_i y_j z_k$$

$$\exp\left(\frac{2\pi i x y z}{2^n}\right) = \prod_{i,j,k} \exp\left(\frac{2\pi i 2^{i+j+k}}{2^n} \mathsf{x}_i \mathsf{y}_j \mathsf{z}_k\right)$$

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 x_i, y_j, z_k are binary values—apply phase only if they all are equal to 1!

A series of CCR_{ϕ} gates between the bits of $|x\rangle$, $|y\rangle$, and $|z\rangle$!

$$\exp\left(\frac{2\pi i xyz}{2^n}\right) = \prod_{i,j,k} \exp\left(\frac{2\pi i 2^{i+j+k}}{2^n} x_i y_j z_k\right)$$

The downside:

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The downside: For n-bit numbers, this requires n^3 gates!

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A modest improvement: classical-quantum multiplication $|\mathcal{U}(a)|x\rangle |0\rangle = |x\rangle |ax\rangle$

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$$\exp\left(\frac{2\pi iaxz}{2^n}\right) = \prod_{i,j} \exp\left(\frac{2\pi ia2^{i+j}}{2^n}x_iz_j\right)$$

Here: $\mathcal{O}(n^2)$ controlled phase rotations (matches Schoolbook algorithm)

Fast quantum multiplication

Main question: Can we combine fast multiplication with Fourier arithmetic to get the benefits of both?

Goal:
$$U(a)|x\rangle|0\rangle = |x\rangle|ax\rangle$$

Goal: Apply phase $\exp\left(\frac{2\pi ia}{2^n}xz\right)$; x and z are quantum

Goal: Implement PhaseProduct
$$(\phi) |x\rangle |z\rangle = \exp{(i\phi xz)} |x\rangle |z\rangle$$

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We want to split the phase ϕxz into the sum of many phases, which are easy to implement.

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Previously:

$$\exp(i\phi xz) = \prod_{i,j} \exp\left(i\phi 2^{i+j} x_i z_j\right)$$

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Karatsuba:

$$xz = 2^{n}x_{1}z_{1} + 2^{n/2}((x_{0} + x_{1})(z_{0} + z_{1}) - x_{0}z_{0} - x_{1}z_{1}) + x_{0}z_{0}$$

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Plugging in Karatsuba:

$$\begin{split} \exp{(i\phi xz)} &= \exp{(i\phi 2^n x_1 z_1)} \\ & \cdot \exp{(i\phi x_0 z_0)} \\ & \cdot \exp{\left(i\phi 2^{n/2} ((x_0 + x_1)(z_0 + z_1) - x_0 z_0 - x_1 z_1)\right)} \end{split}$$

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How are we supposed to reuse values in the phase?

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Re-ordering Karatsuba:

$$xz = (2^{n} - 2^{n/2})x_1z_1 + 2^{n/2}(x_0 + x_1)(z_0 + z_1) + (1 - 2^{n/2})x_0z_0$$

Goal: Implement PhaseProduct
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We want to split the phase ϕ xz into the sum of many phases, which are easy to implement.

Plugging in reordered Karatsuba:

$$\exp(i\phi xz) = \exp\left(i\phi(2^{n} - 2^{n/2})x_{1}z_{1}\right)$$

$$\cdot \exp\left(i\phi(1 - 2^{n/2})x_{0}z_{0}\right)$$

$$\cdot \exp\left(i\phi 2^{n/2}(x_{0} + x_{1})(z_{0} + z_{1})\right)$$

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Goal: Implement PhaseProduct(
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Each of these has the same structure, but on half as many qubits \rightarrow do it recursively!

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$$\exp(i\phi xz) = \exp(i\phi_1 x_1 z_1) \qquad \phi_1 = (2^n - 2^{n/2})\phi$$

$$\cdot \exp(i\phi_2 x_0 z_0) \qquad \phi_2 = (1 - 2^{n/2})\phi$$

$$\cdot \exp(i\phi_3 (x_0 + x_1)(z_0 + z_1)) \qquad \phi_3 = 2^{n/2}\phi$$

Complexity:
$$T(n) = 3T(n/2)$$

Fast classical-quantum multiplication

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Complexity:
$$T(n) = 3T(n/2) \Rightarrow \mathcal{O}(n^{\log_2 3}) = \mathcal{O}(n^{1.58\cdots})$$
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Goal: Apply phase $\exp\left(\frac{2\pi i}{2^n}xyz\right)$; x, y, and z are quantum

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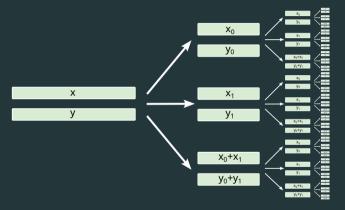
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k	Runtime
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These runtimes are achieved with 2 ancilla qubits.

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Claim: Can implement PhaseProduct and PhaseTripleProduct in sub-linear time, using $\mathcal{O}(n)$ ancillas!

Suprisingly, multiplication itself is bottlenecked by QFT

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Output register requires $n + \mathcal{O}(\log(1/\epsilon))$ qubits

Summary

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1 ancilla qubit

k	Gates
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3	$\mathcal{O}(n^{1.46\cdots})$
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2 ancilla qubits

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Exact QFT in $\mathcal{O}(n^{1.46})$ gates using 1 ancilla

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- · How well can we optimize explicit circuits?

Thank you!

Greg Kahanamoku-Meyer — gkm@berkeley.edu

Backup

Fast classical-quantum multiplication: algorithm

 $\mathsf{PhaseProdu}\overline{\mathsf{ct}(\phi,\ket{x},\ket{z})}$

Input: Quantum state $|x\rangle |z\rangle$, classical value ϕ

Output: Quantum state $\exp(i\phi xz)|x\rangle|z\rangle$

- 1. Split $|x\rangle$ and $|z\rangle$ in half, as $|x_1\rangle$ $|x_0\rangle$ and $|z_1\rangle$ $|z_0\rangle$
- 2. Apply PhaseProduct $((2^n-2^{n/2})\phi,|x_1\rangle\,,|z_1\rangle)$
- 3. Apply PhaseProduct $((1-2^{n/2})\phi,|x_0\rangle,|z_0\rangle)$
- 4. Add $|x_1\rangle$ to $|x_0\rangle$, and $|z_1\rangle$ to $|z_0\rangle$. Registers now hold $|x_1\rangle$ $|x_0+x_1\rangle$ $|z_1\rangle$ $|z_0+z_1\rangle$.
- 5. Apply PhaseProduct $(2^{n/2}\phi, |x_0 + x_1\rangle, |z_0 + z_1\rangle)$.
- 6. Subtract $|x_1\rangle$, $|z_1\rangle$ to return to registers to $|x_1\rangle$ $|x_0\rangle$ $|z_1\rangle$ $|z_0\rangle$.